An infrared study of the L1551 star formation region

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An infrared study of the L1551 star formation region


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Abstract. Spectroscopic observations using the Infrared Space Observatory are reported towards the two well known infrared sources and young stellar objects L1551 IRS 5 and L1551 NE, and at a number of locations in the molecular outflow. The ISO spectrum contains several weak gas-phase lines of OI, CII, [Fe II] and [Si II], along with solid state absorption lines of CO, CO2, H2O, CH4 and CH3OH. Hubble Space Telescope (HST) images with the NICMOS infrared camera reveal a diffuse conical shaped nebulosity, due to scattered light from the central object, with a jet emanating from L1551 IRS 5. The continuum spectral energy distribution has been modelled using a 2D radiative transfer model, and fitted for a central source luminosity of 45 L⊙, surrounding a dense torus extending to a distance of ~3×10^2 AU, which has a total mass of ~13 M⊙. The visual extinction along the outflow is estimated to be ~10 and the mid-plane optical depth to L1551 IRS 5 to be ~120.

This model provides a good fit to the ISO spectral data, as well as to the spatial structures visible on archival HST/NICMOS data, mid-IR maps and sub-millimetre radio interferometry, and to ground-based photometry obtained with a range of different aperture sizes. On the basis of the above model, the extinction curve shows that emission at wavelengths shorter than ~2 μm is due to scattered light from close to L1551 IRS 5, while at wavelengths >4 μm, is seen through the full extinguishing column towards the central source. Several [Fe II] lines were detected in the SWS spectrum towards L1551 IRS 5. Although it would seem at first sight that shocks would be the most likely source of excitation for the [Fe II] in a known shocked region such as this, the line intensities do not fit the predictions of simple shock models. An alternative explanation has been examined where the [Fe II] gas is excited in hot (~4000 K) and dense (~10^9 cm^-3) material located close to the root of the outflow. The SWS observations did not detect any emission from rotationally excited H2. Observations with United Kingdom Infrared Telescope (UKIRT) of the vibrationally excited S- and Q-branch lines were however consistent with the gas having an excitation temperature of ~2500 K. There was no evidence of lower temperature (~500 K) H2 gas which might be visible in the rotational lines. Observations with UKIRT of the CO absorption bands close to 2.4 μm are best fit with gas temperatures ~2500 K, and a column density ~6×10^20 cm^-2.

There is strong circumstantial evidence for the presence of dense (coronal and higher densities) and hot gas (at least 2500 K up to perhaps 5000 K) close to the protostar. However there is no obvious physical interpretation fitting all the data which can explain this.

Key words: ISM: dust, extinction – ISM: individual objects: L1551 – ISM: jets and outflows – stars: pre-main sequence – infrared: ISM: lines and bands – infrared: stars

1. Introduction

Lynds 1551 is one of the most intensively studied molecular outflow sources. Lying at a distance of ~150 pc in the Taurus-Auriga dark cloud, it is associated with a 30 L⊙ Class I protostar, L1551 IRS 5. This is presumed to be in a pre-T Tauri phase and the driving source of a molecular outflow (Snell et al. 1980; Rainey et al. 1987; Fridlund & White 1989a,b; White et al. 1991), and an optical jet (Mundt & Fried 1983; Fridlund & Liseau 1988). L 1551 IRS5 has long been believed to represent...
arche-typically the very early stages of low mass star formation. It is special in that it is closest to the Sun and that it displays its outflow at only a slight inclination angle. In addition, IRS5 is relatively isolated, thus largely reducing source confusion and signal contamination problems, which has made possible observational studies at high resolution and accuracy not obtainable elsewhere. These high resolution studies have revealed that protostars are seemingly far more complex systems than commonly believed. The extinction, $A_v$, towards IRS 5 has been estimated to be $\gtrsim 150$ $^{\text{mm}}$ (Smith et al. 1987; Stocke et al. 1988). Continuum maps reveal that the dense central core is surrounded by an extended cloud (Woody et al. 1989; Keene & Masson 1990; Lay et al. 1994). The spectral energy distribution and intensity maps of L1551 IRS 5 have been modelled in detail using radiative transfer methods in spherical (Butner et al. 1991, 1994) and axially-symmetric (Men’shchikov & Henning 1997, hereafter MH97) geometries. These suggest that a flat accretion disc (Butner et al. 1994) or a geometrically thick torus (MH97) lies inside the extended cloud.

High resolution radio observations have shown evidence for a double source located at IRS 5 (Rodríguez et al. 1986, 1999; Campbell et al. 1988; Bieging & Cohen 1985). Other interpretations of the available data suggest a different morphology, with a binary system lying at the centre of IRS 5 whose components are separated by $\sim 50$ AU (Looney et al. 1997), which is in turn surrounded by a dusty disc. Hubble Space Telescope (HST) observations by Fridlund & Liseau (1988) suggest that there are two distinct optical jets, supporting this circumbinary interpretation, and that the central region is surrounded by a torus, with a mass $\sim 0.1-0.3 M_{\odot}$, and a radius of $\sim 700$ AU. This in turn is surrounded by an $\sim 70$ AU central cavity, which contains the double radio source. There appears to be an evacuated cavity in the torus, with a half-opening angle of about 50–55°. The axis of the molecular outflow is inclined at about 30–35° to the line of sight.

In this paper, we report spectroscopic observations obtained with the ISO Long and Short Wavelength spectrometers (LWS, SWS) towards IRS 5, L1551 NE, HH 29 and at a number of locations along the molecular outflow. Archival data from the HST/NICMOS camera and using an infrared spectrometer at UKIRT are also examined to provide constraints to the modelling. We then present a detailed self-consistent two-dimensional continuum radiative transfer model for IRS 5 which is consistent with the available data.

### 2. The observations

In this study, the Infrared Space Observatory (ISO) spectrometers (ISO – Long wavelength Spectrometer – LWS – Clegg et al. 1996, and Short Wavelength Spectrometer – SWS – de Graauw et al. 1996) were used to obtain spectra towards IRS 5 and at 10 positions on the outflow, as summarised in Table 1. Some of these pointing positions lay along the centre of the outflow axis, and some at the edge of the outflow cavity (see Fig. 1). An ISO SWS spectrum was obtained towards IRS 5 using the S01 mode (2.4–45 $^{\mu}$m, scan speed 4, resolution $\sim 1000–2000$, integration time 6590 seconds). The SWS aperture varied from $14''\times20''$ to $17''\times40''$ for the 2–40 $^{\mu}$m regions respectively (de Graauw et al. 1996). The SWS and LWS data were reduced using the standard ISO analysis software ISAP v1.6 and LIA 6. Small corrections were made for fringing, and to align adjacent detector scans, but overall the standard pipeline data was of high quality. The SWS observations were made with the long dimension of the slit oriented at position angle $171°$ (measured anticlockwise from north) – this lies almost orthogonal to the direction of the outflow jet (see Fig. 2) and limb-brightened outflow lobes (Rainey et al. 1987).

Images were taken from the HST archive to examine the near-infrared distribution closer to IRS 5. These data, using the HST/NICMOS camera were cleaned of residual cosmic rays which remained after the standard pipeline processing and mosaicked together. Narrow band images of the $\lambda 2.12\ H_2$, $\lambda 1.87$ Paschen lines were examined, as well as broad filters which included the $\lambda 2.12\ H_2$ + CO band-head region, the $\lambda 1.64$ [Fe II] line, and a 1.12 $^{\mu}$m continuum filter (in this paper we will adopt

### Table 1. ISO observation log. The fluxes of the O I and C II lines are in Wcm$^{-2}$. All errors and upper limits are 1 $\sigma$

<table>
<thead>
<tr>
<th>Name</th>
<th>RA (2000) h m s</th>
<th>Dec (2000) $\circ$ $'$ $''$</th>
<th>Integration sec</th>
<th>Instrument</th>
<th>O$_{\lambda 63\mu m}$</th>
<th>O$_{\lambda 145\mu m}$</th>
<th>C$_{\lambda 157\mu m}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>IRS 5</td>
<td>4 31 34.3</td>
<td>18 08 05.1</td>
<td>4265</td>
<td>LWS</td>
<td>$1.1 \pm 0.1 \times 10^{-18}$</td>
<td>$\leq 8.5 \times 10^{-21}$</td>
<td>$1.5 \pm 0.3 \times 10^{-19}$</td>
</tr>
<tr>
<td>IRS 5</td>
<td>4 31 34.3</td>
<td>18 08 05.1</td>
<td>6590</td>
<td>SWS</td>
<td>–</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>HH 29</td>
<td>4 31 28.2</td>
<td>18 06 12.7</td>
<td>7133</td>
<td>LWS</td>
<td>$1.5 \pm 0.1 \times 10^{-19}$</td>
<td>$8.6 \pm 3.1 \times 10^{-21}$</td>
<td>$8.3 \pm 0.6 \times 10^{-20}$</td>
</tr>
<tr>
<td>Flow 3</td>
<td>4 31 51.7</td>
<td>17 59 00.1</td>
<td>7133</td>
<td>LWS</td>
<td>$5.9 \pm 1.3 \times 10^{-20}$</td>
<td>$7.0 \pm 4.8 \times 10^{-21}$</td>
<td>$7.8 \pm 0.3 \times 10^{-20}$</td>
</tr>
<tr>
<td>b1</td>
<td>4 31 26.3</td>
<td>18 07 25.0</td>
<td>1937</td>
<td>LWS</td>
<td>$3.8 \pm 1.1 \times 10^{-20}$</td>
<td>$\leq 4.0 \times 10^{-21}$</td>
<td>$9.5 \pm 0.8 \times 10^{-20}$</td>
</tr>
<tr>
<td>b2</td>
<td>4 31 20.5</td>
<td>18 07 42.4</td>
<td>1777</td>
<td>LWS</td>
<td>$1.4 \pm 0.6 \times 10^{-19}$</td>
<td>$6.0 \pm 2.9 \times 10^{-21}$</td>
<td>$8.8 \pm 1.2 \times 10^{-20}$</td>
</tr>
<tr>
<td>b3</td>
<td>4 31 16.4</td>
<td>18 06 45.1</td>
<td>1777</td>
<td>LWS</td>
<td>$2.4 \pm 0.8 \times 10^{-19}$</td>
<td>$\leq 5.1 \times 10^{-20}$</td>
<td>$1.3 \pm 0.1 \times 10^{-19}$</td>
</tr>
<tr>
<td>b4</td>
<td>4 31 12.5</td>
<td>18 05 35.8</td>
<td>1777</td>
<td>LWS</td>
<td>$\leq 1.0 \times 10^{-19}$</td>
<td>$\leq 5.0 \times 10^{-20}$</td>
<td>$7.3 \pm 0.8 \times 10^{-20}$</td>
</tr>
<tr>
<td>r1</td>
<td>4 31 39.8</td>
<td>18 07 27.3</td>
<td>1937</td>
<td>LWS</td>
<td>$\leq 1.8 \times 10^{-20}$</td>
<td>$\leq 3.1 \times 10^{-20}$</td>
<td>$6.7 \pm 0.9 \times 10^{-20}$</td>
</tr>
<tr>
<td>r2</td>
<td>4 31 39.0</td>
<td>18 09 57.4</td>
<td>1777</td>
<td>LWS</td>
<td>$1.7 \pm 0.6 \times 10^{-20}$</td>
<td>$\leq 1.2 \times 10^{-20}$</td>
<td>$7.6 \pm 0.8 \times 10^{-20}$</td>
</tr>
<tr>
<td>L1551 NE</td>
<td>4 31 44.4</td>
<td>18 08 32.4</td>
<td>1777</td>
<td>LWS</td>
<td>$4.3 \pm 0.7 \times 10^{-19}$</td>
<td>$\leq 5.4 \times 10^{-21}$</td>
<td>$1.4 \pm 0.2 \times 10^{-19}$</td>
</tr>
<tr>
<td>r4</td>
<td>4 31 42.2</td>
<td>18 11 24.1</td>
<td>1777</td>
<td>LWS</td>
<td>$\leq 1.0 \times 10^{-19}$</td>
<td>$\leq 7.1 \times 10^{-21}$</td>
<td>$7.7 \pm 1.2 \times 10^{-20}$</td>
</tr>
</tbody>
</table>
which was 1″2 wide. Spectra covering the interval 1.86 to 2.52 μm were reduced using the standard STARLINK package CGS4DR, and calibrated against the standard stars BS 1001 and BS 1665.

3. Infrared continuum

3.1. ISO observations

3.1.1. SWS and LWS continuum

The ISO spectrum of L1551 IRS 5 is shown in Fig.3. The shape of the continuum in the LWS spectral region can to first order be fitted by a 50 K blackbody. There is however no reason to believe that the entire central region can be described by blackbody emission at this single mean temperature. In fact, the emission at short wavelengths exceeds the expected blackbody values, indicating that higher temperature dust grains must be present, and that the absorption efficiency of dust grains falls off with wavelengths (see e.g., Fig.9).

There is a discontinuity in the continuum flux level between 40 and 50 μm – where the SWS and LWS spectra join together (see Fig.3 for the error bars). This may be due either to a calibration mismatch between the SWS and LWS spectrometers, or a consequence of the different beamwidths.

It is known that the size of L1551 IRS 5 at 50 μm is less than 6″ FWHM (Butner et al. 1991), and so it should appear point-like to both the SWS and LWS. One should keep in mind, however, the fact that the beam size of the KAO observations was 14″, much smaller than the apertures used in this study. Since larger beams probe different volumes of the circumstellar envelope, the observations cannot be compared directly, unless a detailed radiative transfer model for the object is developed. In fact, our detailed model of L1551 IRS 5 presented later in this section, shows that this continuum level break could be entirely due to the difference in beam sizes for SWS and LWS spectrometers – although of course there is still systematic uncertainty between the calibrations of the LWS and SWS as discussed earlier.

3.1.2. Solid state features

The ISO spectra are shown in Fig.4 and 5. These include a number of features which are attributable to solid state carriers, as well as several gas phase lines.

The spectral region from 2.8 to 3.8 μm has a broad absorption feature extending across it, made up from contributions from ices at 3.1 μm, O–H stretching band of H₂O from 3 to 3.4 μm, and solid CH₃OH at 3.54, 3.84 and 3.94 μm (Dartois et al. 1999), which sits on the long wavelength side of the H₂O line.

The column density, $N$, of an absorbing species producing a feature of peak optical depth $\tau_{\text{max}}$ and width $\delta \nu$ is given by

$$N = \frac{\tau_{\text{max}} \delta \nu}{A}$$

where $A$ is the integrated absorption cross section per molecule (also known as the band strength) and $\delta \nu$ is the full-width at
half maximum intensity, expressed in units of wavenumbers. The lower limit to the column density of solid H$_2$O is $7.7 \times 10^{18}$ cm$^{-2}$ (the lower limit is a consequence of the available signal-to-noise ratio (S/N), baseline uncertainty and blending with the nearby CH$_3$OH emission). Examples showing the optical depth in other solid state features are presented in Fig. 5.

It proved difficult to unambiguously separate the contributions of CH$_3$OH from the O–H stretching band of H$_2$O, however from first order Gaussian fitting, we estimate the solid CH$_3$OH column density is $2.6 \times 10^{18}$ cm$^{-2}$.

The CO$_2$ $\lambda$ 4.27 stretching mode and CO$_2$ $\lambda$ 15.2 bending mode lines are clearly seen in absorption with optical depths $\sim$ 2 and 1 respectively. The CO$_2$ $\lambda$ 15.2 line has an integrated line flux of $-3.72 \pm 0.5 \times 10^{-18}$ W cm$^{-2}$ and an optical depth of 0.4. Taking values of band strengths from Gerakines et al. (1995), the column densities in these CO$_2$ stretching and bending modes are $5.5 \times 10^{17}$ cm$^{-2}$ and $5.3 \times 10^{17}$ cm$^{-2}$ respectively.

The 'XCN' $\lambda$ 4.619 line reported by Tegler et al. (1993) was not seen in our data, despite detection of the nearby CO$_2$ $\lambda$ 4.38 line. Similarly, the $\lambda$ 4.68 line detected by Pendleton et al. (1999) was also not convincingly present towards L1551 IRS 5.

An absorption feature is seen close to 6 $\mu$m (see also Fig. 5), which is probably due to a combination of water ice, HCOOH and PAH features. Absorption at these wavelengths has previously been proposed by Tielens et al. (1984) and identified by Keane et al. (1999) towards several sources, and seen towards several of the infrared sources towards the Galactic Centre by Chiar et al. (2000), where it is suggested to indicate the presence of a mixture of H$_2$O, NH$_3$, HCOOH and an aromatic C–C component.

An absorption feature associated with the 7.66 $\mu$m CH$_4$ $\nu_4$ 'deformation mode' line is clearly visible, with an optical depth $\tau = 0.07$. Using a band strength of $A = 7.3 \times 10^{-18}$ cm molecule$^{-1}$ (Boogert et al. 1996), we estimate the column density of solid methane to be $7.6 \times 10^{16}$ cm$^{-2}$. Boogert et al. (1996) suggest that the CH$_4$ is embedded with a mixture of polar molecules (such as H$_2$O, CH$_3$OH, CH$_4$ or other species) in the icy grain mantle. However, despite the availability of models of CH$_4$ in various mixtures of polar molecules, the S/N of the present data is insufficient to discriminate between various mixtures.

The column densities estimated above for the various solid state features are summarised in Table 2.
Fig. 4. Lines detected in the L1551 IRS 5 spectrum. The H$_2$O lines indicated in the lower right pane show the expected wavelengths, but do not necessarily indicate detections.

3.1.3. Spectral shape

Compared to many pre-main sequence objects with outflows which have been observed with the ISO LWS, it is striking how few lines have been detected in these very deep spectra (Saraceno et al. 1999). The low intensity of the CII line and the relative invariance of its strength with distance from IRS 5 shows that the UV field must be weak, even close to IRS 5 and the HH 29 region. The non-detection of high-$J$ CO lines suggests that the outflow gas is neither very hot ($T_{	ext{ex}} \lesssim 200$ K), nor very dense ($n_{\text{H}_2} \lesssim 10^4$ cm$^{-3}$).
Fig. 5a–f. Optical depth in various lines, superimposed by model fits to various gas phase lines. These models are run for a gas temperature of 200 K and Doppler width $\nu_D = 2 \text{ km s}^{-1}$, and column densities of $a \ 10^{21} \text{ cm}^{-2}$, $b \ \times 10^{19} \text{ cm}^{-2}$, $c \ \times 10^{20} \text{ cm}^{-2}$ (uppermost model curve at the top of the figure was run for a temperature of 200 K, the lower model curve was for the higher temperature of 2500 K, and shows the bandhead more clearly), $d \ 10^{21} \text{ cm}^{-2}$, $e \ \times 10^{16} \text{ cm}^{-2}$, $f \ \times 10^{19} \text{ cm}^{-2}$. All data and models have an effective resolution of 700.

Table 2. Summary of solid state column densities

<table>
<thead>
<tr>
<th>Species</th>
<th>$\lambda$ (µm)</th>
<th>Band strength (cm$^{-1}$ mol$^{-1}$)</th>
<th>$N_{col}$ (cm$^{-2}$)</th>
<th>$N_X/N_{H_2O}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>H$_2$O</td>
<td>$3$</td>
<td>$2 \times 10^{-16}$</td>
<td>$\geq 7.7 \times 10^{18}$</td>
<td>$1.0$</td>
</tr>
<tr>
<td>CO$_2$</td>
<td>$4.27$</td>
<td>$7.6 \times 10^{-17}$</td>
<td>$5.5 \times 10^{17}$</td>
<td>$0.07$</td>
</tr>
<tr>
<td>CO$_2$</td>
<td>$15.2$</td>
<td>$1.1 \times 10^{-17}$</td>
<td>$5.3 \times 10^{17}$</td>
<td>$0.07$</td>
</tr>
<tr>
<td>CH$_4$</td>
<td>$7.66$</td>
<td>$7.3 \times 10^{-18}$</td>
<td>$7.6 \times 10^{16}$</td>
<td>$0.01$</td>
</tr>
<tr>
<td>CH$_3$OH</td>
<td>$3.54$</td>
<td>$2 \times 10^{-16}$</td>
<td>$\geq 2.6 \times 10^{18}$</td>
<td>$0.3$</td>
</tr>
</tbody>
</table>

3.2. HST NICMOS images

A composite showing images taken from the HST archive is shown in Fig. 6.

The spatial distribution of radiation (Fig. 6) shows a nebular structure with an opening angle of $\sim 60^\circ$. In two of the filters – the [Fe II] and 1.12 µm broad filters, emission from the 'jet' is clearly visible – most prominently in the [Fe II] filter (this contains the $\lambda 1.64$ [Fe II] line). It is notable that three [Fe II] lines are also seen in the ISO SWS spectra (see Sect. IV), and that the 'jet' is not detectable in the narrowband $\lambda 2.12$ H$_2$ filter.

To attempt confirm the lack of H$_2$ emission from the jet, a first order continuum subtraction, the $\lambda 2.12$ H$_2$ + CO broad filter was scaled and subtracted from the narrowband $\lambda 2.12$ H$_2$ filter. The 'continuum subtracted image' formed in this way showed no evidence of H$_2$ emission from the jet down to a level of 2% of the peak $\lambda 2.12$ emission. This may be due to the fact that either there is no detectable H$_2$ emission from the jet, or that the H$_2$ emission is exactly cancelled out by the CO bandhead emission contained in the 2.12 H$_2$ + CO broad filter. Resolution of this will require sensitive echelle spectroscopy which is not yet available. By contrast, [Fe II] emission is clearly visible without any attempt to 'continuum subtract'.

The mass of the torus was estimated to be $0.1–0.3 \ M_\odot$, the full opening angle was $\sim 100–110^\circ$, and the radius 630 AU (Lucas & Roche 1996). A similar bipolar geometry of L1551 IRS 5 has been found also by MH97 in extensive modelling which, among other things, has shown that the actual size and mass of the dusty surroundings of IRS 5 is at least an order of magnitude larger ($\sim 10^4$ AU and $\sim 10 \ M_\odot$, respectively). The distribution of near-IR emission has been studied towards this source by Lucas & Roche (1996), who used image sharpening techniques to obtain a deconvolved angular resolution with a FWHM of 0'.3.
Fig. 6. HST images of IRS 5 obtained with the NICMOS 2 camera (total field of view 19.′2 × 19.′2, pixel size 0.′075). The area shown here has dimensions of 8.′′2 × 8.′′2, and the diffraction limited resolution is ∼ 0.′′18. The peak intensity on the continuum subtracted images is 2 % of the peak flux of the direct 2.12 µm image.

Based on Monte Carlo modelling, they suggested that the light originated from the scattered light associated with a circumstellar torus with an evacuated bipolar cavity.

3.3. Two-dimensional radiative transfer model

In order to interpret the observations in a quantitative way, we constructed a self-consistent two-dimensional (2D) radiative transfer model for L1551 IRS 5. Whereas MH97 have already presented a comprehensive model for this object, their calculations were affected by numerical energy conservation problems resulting from very high optical depths of the model and incomplete convergence of the iterations. The problem, which mainly affected the total luminosity of the central object and the near- to mid-IR part of the SED in the MH97 model, has now been improved (see, e.g., the model of HL Tau by Men’shchikov et al. 1999, hereafter MHF99).

In this paper, we have recomputed the model using the modified version of the code and the new constraints provided by the ISO and HST observations presented above. Our approach and the model are basically the same as those in MH97 and MHF99. We refer to the papers for more detailed discussion of our approach, computational method, model parameters, and ‘error bars’ of the modelling.

3.3.1. Geometry

Following MH97, we assume that the central star (or a binary) is surrounded by a dense core (with a radius of ∼ 100 AU), which is embedded within a much larger non-spherical envelope (outer radius of ∼ 3 × 10^4 AU). A conical cavity has been excavated by the bipolar outflow, and has a full opening angle of 90°. This axially-symmetric geometry is the same for both the core and the surrounding material. The geometry is shown in Fig. 7.

The density distributions inside the torus and in the bipolar cavities are functions of only the radial distance \( r \) from the centre, where the source of energy is located. We neglect in this model the putative binary system inside the dense core, because its semi-major axis (∼ 20 AU) would be much smaller than the radius of the core. If the binary does exist, it is unlikely that there is a very large cavity around it, with a radius of ∼ 20 AU. Our modelling has shown that in the presence of such a dust-free cavity, most of the inner dust boundary would have temperatures of only ∼ 150 K, far too low to explain the observed SED of L1551 IRS 5. In fact, the near- and mid-IR fluxes would be (many) orders of magnitude less than the observed ones. Instead of assuming that the entire binary fits into the dust-free cavity, we adopt the view that a substantial amount of gas and dust exists deeper inside the core, as close as ∼ 0.2 AU to the central source(s) of energy (see Sect. 4.1.1 later).
3.3.2. Grain properties

Assuming a similarity between HL Tau and L1551 IRS 5, we adopted for the latter the dust properties proposed by MHF99. The only difference is that magnesium-iron oxide grains ($\text{Mg}_0.6\text{Fe}_0.4\text{O}$) are absent, because there is no evidence of an emission/absorption signature close to 18 $\mu$m that would warrant introducing another free parameter to into the model (Fig. 5).

1. The large dust particles have an unspecified composition, radii 100–6000 $\mu$m, size distribution exponent $p = -4.2$, dust-to-gas mass ratio $\rho_d/\rho = 0.01$, average material bulk density $\rho_{gr} = 2.0$ g cm$^{-3}$, sublimation temperature $T_{sub} \sim 1700$ K. They show a gray (i.e., independent of wavelength) extinction efficiency for $\lambda < 600$ $\mu$m, whereas at longer wavelengths it falls off as $\lambda^{-0.5}$.

2. Core-mantle grains are assumed to contain silicate cores ($\text{Mg}_0.6\text{Fe}_0.4\text{SiO}_3$) with $\rho_{gr} = 3.2$ g cm$^{-3}$, covered by dirty ice mantles. The ratio of the total core-mantle grain radii to those of the pyroxene cores is 1.4, and their total radii are 0.11–0.7 $\mu$m, $p = -4.2$, $\rho_d/\rho = 0.0037$. The dirty mantles consist of water ice polluted by small amorphous carbon grains. The sublimation temperature of the mantles was assumed to be $\sim 100$ K.

3. Amorphous carbon grains with radii 0.08–0.5 $\mu$m, $p = -4.2$, $\rho_d/\rho = 0.0063$, $\rho_{gr} = 2.0$ g cm$^{-3}$.

The contributions to the extinction towards L1551 IRS 5 are shown in Fig. 8.

Fig. 7. Geometry of L1551 IRS 5. Schematically shown are three regions of the model – the innermost dense torus (dark color), the constant-density part of the envelope (medium color), and the outer extended envelope containing most of the mass (light color). The bipolar geometry is defined by the opening angle of the conical outflow cavities, $\omega = \pi - \psi \approx 90^\circ$ ($\psi \approx 90^\circ$) and the viewing angle, $\theta_{\psi} \approx 44^\circ$, between the equatorial plane and the line of sight. The polar outflow regions are less dense than the torus.

The first component of very large grains is present only inside the dense torus ($0.2 \text{ AU} \leq r \leq 250 \text{ AU}$), where all of the smaller grains are assumed to have grown into the large particles. The two other grain components exist only outside ($250 \text{ AU} \leq r \leq 3 \times 10^4 \text{ AU}$), in the extended envelope of much lower density. As in any model, the results depend on the accuracy of the input parameters and the assumptions made. A major source of uncertainty in a model such as the one used here, will be the assumptions adopted for the grain properties. Further details on the choice of grain properties and their effect on the model results have been discussed by MHF99 and references therein.

3.3.3. Parameter space

As has been discussed in detail by MH97 and MHF99, the parameter space available to numerical modelling is very large. There are many poorly constrained parameters, few of which can be fixed a priori, to reduce the space. This situation requires that all available observational information has to be taken into account to better constrain the models. As in the earlier modelling, we have used as observational constraints all existing photometry data with different beam sizes (from optical to millimetre wavelengths), intensity profiles at 50 $\mu$m, 100 $\mu$m, 1.25 mm, 1.3 mm, and visibility curves from interferometry at 0.87 mm and 2.73 mm (see references in MH97). In addition to the constraints, we used in the new modelling our HST image at 2.12 $\mu$m and the SWS and LWS spectrophotometry presented above.
Numerical parameters related to the accuracy of the model are the number of radial points (277), the number of azimuthal angles for the integration of intensity moments (10), the number of azimuthal angles for the observable flux calculation (50), the number of wavelengths (217), the number of points for convolved intensity maps (600 × 600), and that for the visibility calculations (4000 × 4000). Conservation of the total luminosity for both the equivalent spherical envelope and the 2D model was better than about 7% at all radial points, certainly good enough compared to the total uncertainties involved in the modelling.

Table 3 lists the main input parameters of our model. The opening angle has not been varied, being fixed at the value found by MH97. Likewise, the viewing angle and the outer boundary also have not been varied extensively in our new modelling. We varied mainly the radial density profile and the total mass (or the total optical depth in the mid-plane). The dust grain parameters, the primary source of uncertainties in this kind of modelling, was also mostly fixed at the values adopted by MH97 for HL Tau.

Numerical parameters related to the accuracy of the model are the number of radial points (277), the number of azimuthal angles for the integration of intensity moments (10), the number of azimuthal angles for the observable flux calculation (50), the number of wavelengths (217), the number of points for convolved intensity maps (600 × 600), and that for the visibility calculations (4000 × 4000). Conservation of the total luminosity for both the equivalent spherical envelope and the 2D model was better than about 7% at all radial points, certainly good enough compared to the total uncertainties involved in the modelling.

3.3.4. Spectral energy distribution

The model SED is compared to the observations of L1551 IRS 5 in Figs. 9 and 10. Only one, best-fitting SED corresponding to the viewing angle of 44°5 (measured from the mid-plane) is shown. As has been extensively discussed by MH97 and MHF99 (see, e.g., their Fig. 5), the optical to mid-IR parts of the SEDs of bipolar embedded sources depend strongly on the viewing angle. The viewing angle derived on the basis of 2D radiative transfer calculations depends, among other factors, on the density distribution in the polar direction inside the dense
The overall quantitative agreement of the model SED with the entire set of observations of L1551 IRS 5 is obvious. The total model fluxes corrected for the beam sizes (solid lines in Figs. 9 and 10) coincide well with the observed fluxes, except for those in the near IR, although the shape of the SED is still very similar to the observed one. The effect of different apertures is evident everywhere, except for only the mid-IR wavelengths, where the source is very compact and most of its radiation fits into the SWS beam. Note that at millimetre waves the model predicts significantly larger total fluxes compared to the observed ones, indicating that the outer envelope is very extended and sufficiently massive.

The insert in Fig. 10 shows that the ‘jump’ between the SWS and LWS data at \( \sim 45 \mu m \) is a consequence of the different beam sizes. The model shows a clear water ice absorption feature at \( 3 \mu m \) which is very similar to the observed profile. The agreement of the model with SWS in the \( 7–9 \mu m \) region is not very good; there are also smaller deviations in the \( 15–40 \mu m \) part of SWS. As we mostly fixed the grain properties, we have not attempted to find a better fit by varying the dust model (it would be extremely time-consuming). It seems very likely, however, that small changes in the dust chemical composition or temperature profile would be sufficient to fit the SED almost perfectly. We do not believe, however, that such an adjustment makes sense, given the much higher overall uncertainties of the problem at hand.

3.3.5. Densities and temperatures

The structure of our model of L1551 IRS 5, which is very similar to that presented by MH97, is illustrated in Fig. 11. The distribution of densities and temperatures in the model were chosen to be similar to those of HL Tau (MHF99), except for the flat density area between 250 and 2000 AU which is very likely to exist in IRS 5. The density structure in the inner few thousand AU is constrained by the SED (Sect. 3.3.4), the submm/mm visibilities (Sect. 3.3.6), and the long-wavelength intensity maps (Sect. 3.3.7). The visibilities suggest that the density structure consists of a dense core inside a lower density envelope. On the other hand, the intensity maps might also imply high densities are present at distances of \( \sim 4000 \) AU from the central source. Both requirements can however only be reconciled by adopting a flat density distribution. We have no clear understanding of the physical significance of this deduction (also inferred by MH97), which needs to be tested by other observations.

There are three regions that make up the torus: the innermost very dense core with a \( \rho \propto r^{-1} \) density gradient, and low-density outer parts with a broken power-law \( (\rho = \text{const}, \rho \propto r^{-\gamma}) \) density profile. A steep \( \rho \propto \exp(-r^2) \) transition zone between them (having a half-width at half-maximum of 70 AU) effectively forms the outer boundary of the inner dense torus. The boundary of the torus extends from \( \sim 80 \) to 250 AU and is effectively truncated by the exponential at about 200 AU, very similar to the density profile of HL Tau (MHF99). Conical surfaces of the bipolar outflow cavities define the opening angle of the torus to be \( 90^\circ \). Dust evaporation sets the inner boundary at \( \approx 0.4 \) AU, while the outer boundary is arbitrarily put at a sufficiently large distance of \( 3 \times 10^4 \) AU. The polar outflow cones with a \( \rho \propto r^{-2} \) density distribution have much lower density than the torus.

In the absence of any reliable constraints, the conical outflow regions are assumed to have a \( \rho \propto r^{-2} \) density profile which is consistent with available data. The temperature profile displays a jump at 250 AU, where the hotter normal-sized dust grains of the envelope are assumed to be coagulated into the large grains of the dense core. We refer to MH97 and MHF99 for a more detailed discussion of the density structure and of the uncertainties of our model.

3.3.6. Submillimetre and millimetre visibilities

The model visibilities for two directions in the plane of sky, parallel and orthogonal to the projected axis of the torus (MH97, MHF99) are compared to the available interferometry data in Fig 12. The model shows good agreement with the spatial information contained in the observed visibilities. The latter do not constrain, however, density distribution in the outer envelope. Instead, intensity maps over a larger area, obtained with large
beams, should be useful in determining the density structure on the largest scales, thus giving an idea of the total mass of the circumstellar material. In fact, our model gives much larger mass ($\sim 13 M_\odot$) and extent ($\sim 3 \times 10^4$ AU) of the envelope compared to other simplified models which do not take into account all available observations.

The observations by Looney et al. (1997) were added to the figure after our modelling has been completed; they show noticeably lower visibilities compared to those presented by Keene & Masson (1990), which were fit by the model reasonably well. This is clearly a consequence of the assumed $\lambda^{-0.5}$ slope of the opacities by very large grains in this wavelength range (Sect. 3.3.2). The slope has been chosen in our model on the basis of the opacities by very large grains in this wavelength range. The latter data set was not used in the search through the parameter space of the model (see Sect. 3.3.6). The upper and lower curves show the visibilities for two directions in the plane of sky, parallel and orthogonal to the projected axis of the outflow. The observations by Looney et al. (1997) were added to the figure after our modelling has been completed; they show noticeably lower visibilities compared to those presented by Keene & Masson (1990), which were fit by the model reasonably well. This is clearly a consequence of the assumed $\lambda^{-0.5}$ slope of the opacities by very large grains in this wavelength range (Sect. 3.3.2).

3.3.7. Far-IR and millimetre intensity profiles

Model intensity profiles at 50 $\mu$m, 100 $\mu$m, 1.25 mm, and 1.3 mm (perpendicular to the outflow direction) are compared with the available observations in Figs. 13 and 14. The model intensity profiles have been convolved with the appropriate circular Gaussian beams. Unconvolved intensity distributions are also shown for reference. As in MH97, the new model shows very good agreement with the measurements, suggesting that the density and temperature distributions of the model are realistic.

Note that the temperature of the outermost parts may be controlled by the external radiation field, which is assumed to be a 5 K blackbody in our model. The radiation field defines the lower limit for the dust temperature in the distant parts of the torus. Thus, it is a key parameter that determines how extended the envelope would appear to millimetre observations after subtraction of the background radiation field. In fact, if the external radiation field keeps the torus warm (e.g., $\sim 20$ K), then the millimetre intensity maps, which are sensitive to much cooler material, would reveal very little radiation from the envelope. One cannot conclude, however, that an envelope has very little mass on the basis of a few millimetre fluxes alone, without a careful analysis of its density and temperature structure.

3.3.8. HST image at 2.12 $\mu$m

An additional check of our model at short wavelengths is enabled by the HST NICMOS image of IRS 5 at 2.12 $\mu$m. In Fig. 15, we compare the model intensity profiles along the same orthogonal directions on the sky, parallel and perpendicular to the axis of the outflow, with the corresponding intensity strips taken from the observed images (Sect. 3.2). The model images were convolved with the HST point-spread function of 0
d
3.3.9. Average density estimates

Table 4 compares average densities of our model with the observational estimates. The density reported by Keene & Masson (1990) (see MH97) from 2.7 mm interferometry was an overestimate because they assumed a steep long-wavelength absorption efficiency of dust grains \( \kappa \propto \lambda^{-2} \). Corrected for the higher opacity of large grains adopted in our work, the estimate would be about a factor of 4 times lower, in much better agreement with our result. The model is consistent with observations, given the approximations involved in the problem (see MHF99 for more discussion).

Table 4. Average density estimates for L1551 IRS 5 from observations compared to the predictions of our model (references are given in MH97, Table 2).

<table>
<thead>
<tr>
<th>Tracer</th>
<th>Rad. ( m )</th>
<th>Rad. AU</th>
<th>Obs. ( \langle n_{\text{H}_2} \rangle ) cm(^{-3} )</th>
<th>Mod. ( \langle n_{\text{H}_2} \rangle ) cm(^{-3} )</th>
<th>Ratio</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \text{NH}_3 )</td>
<td>60.0</td>
<td>9600</td>
<td>( \approx 1 \times 10^5 )</td>
<td>( 1.7 \times 10^5 )</td>
<td>( \approx 1.7 )</td>
</tr>
<tr>
<td>1.1 mm</td>
<td>30.0</td>
<td>4800</td>
<td>( \approx 4 \times 10^5 )</td>
<td>( 5.0 \times 10^5 )</td>
<td>( \approx 1.3 )</td>
</tr>
<tr>
<td>400 ( \mu )m</td>
<td>17.5</td>
<td>2800</td>
<td>( \gtrsim 1 \times 10^5 )</td>
<td>( 8.0 \times 10^5 )</td>
<td>( \gtrsim 0.8 )</td>
</tr>
<tr>
<td>( \text{C}^{18}\text{O} )</td>
<td>13.8</td>
<td>2208</td>
<td>( \gtrsim 4 \times 10^5 )</td>
<td>( 8.0 \times 10^5 )</td>
<td>( \gtrsim 2.0 )</td>
</tr>
<tr>
<td>1.4 mm</td>
<td>4.0</td>
<td>640</td>
<td>( \approx 1 \times 10^7 )</td>
<td>( 3.7 \times 10^6 )</td>
<td>( \approx 0.4 )</td>
</tr>
<tr>
<td>( \text{C}^{18}\text{O} )</td>
<td>3.5</td>
<td>560</td>
<td>( \gtrsim 3 \times 10^7 )</td>
<td>( 5.0 \times 10^6 )</td>
<td>( \gtrsim 2.0 )</td>
</tr>
<tr>
<td>2 ( \mu )m</td>
<td>3.0</td>
<td>480</td>
<td>( \gtrsim 4 \times 10^6 )</td>
<td>( 7.0 \times 10^5 )</td>
<td>( \gtrsim 1.8 )</td>
</tr>
<tr>
<td>2 ( \mu )m</td>
<td>2.0</td>
<td>320</td>
<td>( \gtrsim 1 \times 10^6 )</td>
<td>( 2.4 \times 10^5 )</td>
<td>( \gtrsim 2.0 )</td>
</tr>
<tr>
<td>2.7 mm</td>
<td>0.4</td>
<td>64</td>
<td>( \gtrsim 1 \times 10^{11} )</td>
<td>( 1.0 \times 10^9 )</td>
<td>( \leq 0.01 )</td>
</tr>
</tbody>
</table>
4. Emission lines

Several narrow emission lines were detected in the L1551 IRS 5 spectrum, specifically [Fe II] λ 17.94 and λ 25.99, Si II λ 31.48, O I λ 63.2, OH λ 84.4 and C II λ 157.7. These can be seen in Fig. 4 and are now discussed by species.

4.1. [Fe II]

[Fe II] line emission is often detected towards starburst galaxies, where it has been interpreted as having been excited by collisional excitation in supernova remnant shocks (Moorwood & Oliva 1988; Lutz et al. 1998). [Fe II] lines have been reported towards several supernova remnants (Oliva et al. 1999a [RCW 103], b [IC 443]), galactic nuclei (Lutz et al. 1996 [Sgr A*], 1997 [M82]), and in regions known to have energetic outflows (Wesselius et al. 1998 [S106 IR and Cepheus A]). Optical echelle spectra towards the jet and working surface where the jet intersects with the surrounding medium (Fridlund & Liseau 1988) show linewidths of 100 and 200 km s$^{-1}$ respectively. These results unambiguously show that shock excitation is occurring, although the position of the working surface in L1551 IRS 5 lies outside the SWS field of view. From the raw data, the linewidth of the λ 26.0 [Fe II] line, deconvolved from the instrumental resolution, is ≲ 230 km s$^{-1}$. The ratios [Fe II] λ 35.3/λ 26.0 and [Fe II]/λ 25.9 should be density sensitive, although their variation over a wide range of densities is only a factor of ≈ two. The transitions have high critical densities $n_{cr} \sim 10^{10}$ cm$^{-3}$ and high excitation temperatures $T_e \gtrsim 400$ K. As minor coolants, they do not significantly affect the thermal structure of the cloud. To test the possibility that the gas could be photoionised, we ran the photoionisation code CLOUDY (Ferland et al. 1998) over a wide range of values. The [Fe II] λ 24.5 line is only efficiently excited under conditions of high density ($n_{H_2} \gtrsim 10^6$ cm$^{-3}$) and high UV illumination ($G_0 \gtrsim 10^8$) (see also Hollenbach et al. 1991) – however, it is impossible to excite the [Fe II] λ 17.9 line for reasonable values of ionising flux near L1551 IRS 5, due to the low effective temperature of the star ($T_{eff} \sim 5500$ K). A shocked environment seems a more likely environment to excite the lines. Such a shocked environment could be associated with the optical jet emanating from L1551 IRS 5. We detect emission from [Fe II] λ 17.9 and λ 26.0 of $9.44 \pm 1.5 \times 10^{-16}$ W m$^{-2}$ and $2.04 \pm 0.3 \times 10^{-15}$ W m$^{-2}$ respectively, with a marginal detection of [Fe II] λ 35.3 of $1.62 \pm 0.6 \times 10^{-15}$ W m$^{-2}$. There are further [Fe II] lines at λ 5.34, λ 51.3 and λ 87.4 for which we set 2σ upper limits of $1.05 \times 10^{-15}$, $6.52 \times 10^{-16}$ and $3.6 \times 10^{-16}$ W m$^{-2}$ respectively. The λ 5.34 and λ 17.9 lines can be used to constrain density, and the λ 26 line can be used with other lines as a temperature estimator. However, as pointed out by Greenhouse et al. (1997), Lutz et al. (1998) and Justtanont et al. (1999), other excitation mechanisms such as fluorescence or photoionisation may be important in certain environments.

4.1.1. Observed transitions

A diagram of the lowest energy levels of [Fe II] is shown in Fig. 17 and the major transitions in the range of the ISO spectrometers are listed in Table 5. The observed line fluxes are also given, along with the aperture sizes of the data. From these it is clear that for extended emission several of the line intensity ratios are aperture dependent. Fortunately, some of the potentially most important lines, viz. the 26, 17.9 and 24.5 μm lines (the 2–1, 7–6 and 8–7 transitions in Fig. 17 respectively), were observed with the same aperture.

The table also lists intrinsic [Fe II] line fluxes, $F_0$, i.e. after a correction for the attenuation by dust extinction has been applied. The values of the dust opacities used for this correction were obtained from the best fit model of the overall SED of IRS 5 (see Sect. and below) and are displayed in Figs. 8 and 18. Contrary to naive expectations, extinction has a considerable effect on the line ratios in IRS 5 even in the mid- to far-infrared. Not only does the temperature diagnostic ratio 17.9/26 become slightly altered, but is actually inverted, changing from the observed value of 0.5 to the ‘de-reddened’ one of 2.

Emission lines of [Fe II] detected in the SWS spectrum of L1551 IRS 5 include transitions at 26 and 17.9 μm. Further, the high excitation lines (cf. Fig. 17) [Fe II] λ 1.64 (Sect. and λ 0.716 (Fridlund & Liseau, in preparation) are also clearly present in this source, as well as the [Si II] λ 35 line. Therefore, in the absence of any nearby bright source of UV radiation, excitation of these lines by shocks presents the only known, feasible alternative. However, the comparison of the line intensity ratios for various lines deduced from shock models (Hollenbach & McKee 1989; Hollenbach et al. 1989) with those of the observations makes no immediate sense. In fact, results obtained from
within the field of view of the SWS; Fridlund et al. 1997). This is, I
ble 5), would imply an intensity of the non-extinguished line
the [Fe II] emission lines. For instance, the observed flux of
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Fig. 18. Total dust optical depth towards the centre of IRS 5 (in the SWS wavelength range) according to our 2D radiative transfer model (Sect. 3.3). Also shown are the relative contributions from absorption (dotted line) and scattering (dashed line). The positions of the [Fe II] lines are indicated by the vertical bars. Note specifically the extinction bump, due to silicate particles, at the position of the [Fe II] 17.9 µm line.

et al. (1996). The number of radiative transitions included in the calculation was 1438. These lines are distributed from the FUV to the FIR spectral regions (0.16 to 87 µm) and the level energies span the range $E/k = 0 - 9.1 \times 10^4$ K (the ionisation potential of [Fe II] corresponds to nearly $1.9 \times 10^5$ K). The collision rate coefficients were calculated from the work by Zhang & Pradhan (1995), who provide effective collision strengths for 10 011 transitions among 142 fine structure levels in [Fe II]. These are Maxwellian averages for 20 temperatures in the range 1000 K to $10^5$ K.

No lines of [Fe I] or [Fe III] (or of higher ionisation for that matter) have been detected from IRS 5, so that our assumption that essentially all iron is singly ionised seems reasonably justified. We further assume that iron is undepleted in the gas phase with solar chemical abundance, i.e. $A(Fe) = 3.2 \times 10^{-5}$ (Grevesse & Sauval 1999). At temperatures significantly below 8000 K the gas would be only partially (hydrogen) ionised. In this case, we assume that the electrons are donated by abundant species with similar and/or lower ionisation potentials. In particular, primarily by Fe and Si plus other metals such as Mg, Al, Na, Ca etc., so that $n_e \sim 2.5 A(Fe) n(H)$.

The line intensities were calculated for a range of gas kinetic temperatures and hydrogen densities and an example of the results is presented in Fig. 19. In the figure, intensity ratios for [Fe II] 5.34, 17.9 and 26.0 µm lines are shown. These lines are connected (see Fig. 17): the 17.9 µm and 5.34 µm lines originate from the same multiplet, $^4$Fe. The upper level of the 17.9 µm line is at 3500 K above ground and its lower level is the upper level of the 5.34 µm line, which connects to the ground, and so does the 26 µm groundstate line ($^4$De). The corresponding SWS data are also shown in the graph, where the emitting regions have been assumed to be much smaller than any of the apertures used for the observations. These data indicate source temperatures to be somewhere in the range of 3000 to 5000 K and gas densities to be above $3 \times 10^7$ cm$^{-3}$.

4.1.3. [Fe II] model calculations and results

The model spectrum, ‘tuned’ to the SWS observations with the parameters of Sect. 4.1.1, is shown in Fig. 20. The upper panel of the figure displays most of the 1438 lines of the computed intrinsic [Fe II] spectrum, stretching from the far-UV to the far-IR. This spectrum suffers extinction by the circumstellar dust before it reaches the outside observer and is shown in the middle panel. Scattering by the circumstellar dust is of only low significance for the emission in the mid-infrared but is probably important at near-infrared and shorter wavelengths. In the lower panel of Fig. 20, a simple scattering model has been ap-
modelling of the H\textsubscript{2} line radiation about the central source. The dotted horizontal line is meant to aid the eye.

Table 6. [Fe II] line optical depths and intensities (Model: \(T_{\text{kin}} = 4000\,\text{K}, n(H) = 10^9\,\text{cm}^{-3}, \log \frac{X}{(\text{dv}/\text{dr})} = -10.46\))

<table>
<thead>
<tr>
<th>Transition</th>
<th>(\lambda) ((\mu\text{m}))</th>
<th>(\tau_{\text{line}})</th>
<th>(I_{\text{line}}) ((\text{erg cm}^{-2},\text{s}^{-1},\text{sr}^{-1}))</th>
</tr>
</thead>
<tbody>
<tr>
<td>2-1</td>
<td>25.9896</td>
<td>2.26\times10^{-2}</td>
<td>1.96\times10^{4}</td>
</tr>
<tr>
<td>3-2</td>
<td>35.3519</td>
<td>2.05\times10^{-2}</td>
<td>7.22\times10^{4}</td>
</tr>
<tr>
<td>4-3</td>
<td>51.2847</td>
<td>1.20\times10^{-2}</td>
<td>1.42\times10^{4}</td>
</tr>
<tr>
<td>5-4</td>
<td>87.4126</td>
<td>4.34\times10^{-3}</td>
<td>1.05\times10^{-1}</td>
</tr>
<tr>
<td>6-1</td>
<td>5.3402</td>
<td>1.21\times10^{-4}</td>
<td>9.24\times10^{4}</td>
</tr>
<tr>
<td>7-6</td>
<td>17.9379</td>
<td>1.46\times10^{-2}</td>
<td>3.75\times10^{4}</td>
</tr>
<tr>
<td>8-7</td>
<td>24.5182</td>
<td>1.15\times10^{-2}</td>
<td>1.20\times10^{4}</td>
</tr>
<tr>
<td>9-8</td>
<td>35.7731</td>
<td>5.20\times10^{-3}</td>
<td>1.78\times10^{4}</td>
</tr>
<tr>
<td>10-6</td>
<td>1.6440</td>
<td>5.62\times10^{-5}</td>
<td>5.79\times10^{1}</td>
</tr>
<tr>
<td>17-6</td>
<td>0.7157</td>
<td>1.58\times10^{-4}</td>
<td>2.36\times10^{2}</td>
</tr>
<tr>
<td>25-6</td>
<td>0.5160</td>
<td>3.18\times10^{-4}</td>
<td>2.46\times10^{2}</td>
</tr>
</tbody>
</table>

any uncertainty in the computed extinction curve. It is clear that future modelling of this source will need to address a wider parameter space, particularly of grain properties – however this is beyond current computational capabilities.

The spectrum of the adopted model (Table 6) is shown superposed onto selected regions of the observed SWS scans in Fig. 21. The theoretical fit is acceptable for most lines, except perhaps for the 24.52 \(\mu\text{m}\) transition which appears too strong in comparison with the observed line. All of the major [Fe II] lines are thermalised and optically thin, implying that the emission model can be easily scaled by keeping the parameter \(N/\Delta v\) constant. For instance, models with higher velocity and lower density (e.g. 300 km s\(^{-1}\) and \(5\times10^8\) cm\(^{-3}\)) or vice versa (e.g. 15 km s\(^{-1}\) and \(10^{10}\) cm\(^{-3}\)) would still yield the same [Fe II] intensities (but would otherwise disagree with the presence or absence of other lines in the spectrum of IRS 5). The cooling of the gas in all [Fe II] lines amounts to a few times \(10^{-3}\) \(L_\odot\), which is comparable to the [O I] 63 \(\mu\text{m}\) luminosity, \(> 7\times10^{-3}\) \(L_\odot\). For comparison, the total luminosity generated by the shocks can be estimated as \(\max L_n = 0.5 n_s^3 m_H n \times \text{area} \sim 10 L_\odot\), where we have assumed a gas compression \(n/n_0 \sim 10^2\). This amounts to about 20 % of the total radiative luminosity of IRS 5 (~45 \(L_\odot\), Table 3). The largest uncertainty lies in the value of \(n_s\), the dependence of which is cubic. However, the radiative luminosity of the IRS 5 system can be expected to be dominated by accretion processes, whereas the shock luminosity is prob-
ably generated by mass outflows. To order of magnitude, these estimates would then seem reasonable.

As to how the degree of ionisation of the gas, albeit low, is generated and maintained we have essentially no information. Dissociation, if initially molecular, and subsequent ionisation through shocks seems a likely option. In any case, one would expect the partially ionised gas to produce free-free continuum emission, with a flux density (in mJy) $S_\nu = 5.44 \times 10^{-13} g_{\text{f}}(\nu, T) Z^2 T^{-0.5} E M \exp(-h \nu / k T) \Omega$. From the [Fe II] model, the emission measure of the gas is $E M = \int x^2 u^2(H) d s = 6 \times 10^6 \text{cm}^{-6} \text{pc}$, the free-free Gaunt factor $g_{\text{f}}(1.4 \text{GHz}, 4000 \text{K}) = 5$ and the source solid angle $\Omega = 5 \times 10^{-12} \text{sr}$. Consequently, we find at 1.4 GHz (21 cm wavelength) $S_{1.4, \text{obs}} = 3.9 \text{mJy}$, which is not far from what has actually been observed. The Very Large Array (VLA) measurements taken in August 1992 by Giovanardi et al. (2000) obtained $S_{1.4, \text{obs}} = 3.3 \pm 0.3 \text{ mJy}$ are closest in time to the ISO observations.

Recent observations with the Subaru telescope by Itoh et al. (2000) have shown that the optical jet is dominated by [Fe II] lines, and suggest that the extinction to the jet is on average $A_{\text{v}} \sim 7$ mag. Fridlund et al. (1997) provide observed H $\alpha$ fluxes for the entire jet (ground based and HST), viz. $F(\text{H}\alpha)_{\text{obs}} = 4.2 \times 10^{-14}$ erg s$^{-1}$ cm$^{-2}$ (the working surface, knot D, alone radiates 50 erg cm$^{-2}$ s$^{-1}$). The intrinsic line ratio (H $\beta$/H $\alpha$) of $6 \times 10^{-11} - 10^{-9}$ erg s$^{-1}$ cm$^{-2}$. This exceeds by more than one order of magnitude, the value inferred for the jet putatively extinguished by 7 magnitudes of visual extinction.

Since the intrinsic line ratio (H $\alpha$/H $\beta$)$_{\text{intrinsic}}$ is equal to or larger than 3, dust extinction with $A_{\text{v}} = 7$ mag would result in a line ratio (H $\alpha$/H $\beta$)$_{\text{ext}} \geq 35$ (see, e.g. Appendix B of Fridlund et al. 1993). The observed ratio is (H $\alpha$/H $\beta$)$_{\text{obs}} = 15$ (Cohen & Fuller 1985), which is significantly smaller than the predicted lower limit to the line ratio. This would increase the discrepancy even more.

In summary, it is again concluded that the [Fe II] emission observed by the ISO-SWS is not dominated by the jet, but its source is of different origin. That the jet is emitting in [Fe II] lines has been known for some time and is not new, but the level of emission is not sufficient to explain the SWS observations.

4.2. Si II

The sole [Si II] line detected is the $^2P_{3/2} - ^2P_{1/2}$ ground state magnetic dipole transition at $\lambda 34.8$. It should be one of the major coolants in hot ($T \geq 5000$ K) gas. It has been previously suggested to be a shock tracer (Haas et al. 1986). We detect a flux of $2.35 \pm 0.3 \times 10^{-15}$ W m$^{-2}$ from this line. It is of interest to understand to what extent this flux is consistent with the prediction from our model of the central source.

$$F_{\text{line}} = \frac{\Omega_{\text{source}}}{4 \pi} h \nu_{\text{ul}} A_{\text{ul}} n_{\text{ul}} \Delta \ell$$

(2)

with obvious notations. The fractional population of the upper level, $f_u$, is obtained from

$$f_u = \frac{n_{\text{ul}}}{n(\text{Si II})} = \left[1 + \frac{A_{\text{ul}} + n_e q_{\text{ul}}}{n_e q_u}\right]^{-1}$$

(3)

where the collision rate constants, $q(T_e)$ (cm$^3$ s$^{-1}$), are related to the respective Maxwellian average of the collision strength, $\gamma(T_e)$, by

$$q_{\text{ul}} = \frac{8.6287 \times 10^{-6}}{g_{\text{ul}} T_e^{1/2}} \gamma_{\text{ul}}$$

(4)

and

$$q_u = \frac{g_u}{g_{\text{ul}}} q_{\text{ul}} \exp\left(-\frac{h \nu_{\text{ul}}}{k T_e}\right).$$

(5)

At the 2$\sigma$ level, upper limits can be set for $F(\text{Si II}) 68.5 \mu m < 6.4 \times 10^{-20}$ W cm$^{-2}$ and $F(\text{Si III}) 38.2 \mu m < 3.7 \times 10^{-19}$ W cm$^{-2}$, whence we can safely assume that essentially all silicon is singly ionised, $n_{\text{ul}}$ can therefore be expressed as $f_u A(\text{Si}) n(\text{H})$, where $A(\text{Si})$ is the abundance of silicon with respect to hydrogen nuclei.

Energy levels and Einstein-A values were adopted from Wiese et al. (1966) and Kaufman & Sugar (1986) and collision strengths from Callaway (1994). As for iron, silicon is assumed to be undepleted in the central core regions of L 1551 IRS 5 and we adopted the solar value of the silicon abundance, $A(\text{Si}) = 3.2 \times 10^{-5}$ (Asplund 2000).

For the values of the model parameters, viz. $n(\text{H}) = 10^6$ cm$^{-3}$, $n_e = 8 \times 10^4$ cm$^{-3}$ and $T_e = 4000$ K, the fractional population becomes $f_u = 0.64$. Further, in conjunction with $\Omega_{\text{source}} = 8.2 \times 10^{-12} \text{sr}$ and $\Delta \ell = 90$ AU, the intrinsic model flux then becomes $F_0([\text{Si II}] 35 \mu m) = 2.2 \times 10^{-10}$ erg cm$^{-2}$ s$^{-1}$. Taking the dust extinction by the intervening disk into account, $\tau_{35 \mu m} = 4.1$ Fig [18] leads to the predicted estimation of the observable flux, i.e. $F_{\text{model}}([\text{Si II}] 35 \mu m) = 3.6 \times 10^{-12}$ erg cm$^{-2}$ s$^{-1}$.

In this undepleted case, the model flux is only slightly larger, by a factor of $\sim 1.5$, than the actually observed value. We conclude therefore that our model of the central regions in L 1551 IRS 5 is capable of correctly predicting the flux in the [Si II] 35 $\mu$m line. It is likely that most of this emission also originates in these central regions.

4.3. OH

The excitation of OH has been modelled by Melnick et al. (1987), who studied OH emission towards the Orion Nebula.
### Table 7. Summary of IRS 5 H₂ line fluxes

<table>
<thead>
<tr>
<th>Line</th>
<th>λ (μm)</th>
<th>Flux W m⁻²</th>
<th>I_line/I_S(0)</th>
<th>Thermal PDR</th>
<th>High A_v</th>
</tr>
</thead>
<tbody>
<tr>
<td>H₂ S(0)</td>
<td>2.223</td>
<td>7.9 ± 0.8 × 10⁻¹⁸</td>
<td>5.6 ± 0.1 × 10⁻¹⁸</td>
<td>0.72 0.93 0.21</td>
<td>0.07 0.06 0.44</td>
</tr>
<tr>
<td>H₂ S(1)</td>
<td>2.122</td>
<td>1.1 ± 0.1 × 10⁻¹⁷</td>
<td>0.6 ± 0.1 × 10⁻¹⁸</td>
<td>1.00 1.00 1.00</td>
<td>1.00 1.00 1.00</td>
</tr>
<tr>
<td>H₂ S(2)</td>
<td>2.034</td>
<td>3.9 ± 0.4 × 10⁻¹⁸</td>
<td>0.5 ± 0.1 × 10⁻¹⁸</td>
<td>0.36 0.37 0.28</td>
<td>0.13 1.59 0.23</td>
</tr>
<tr>
<td>H₂ S(3)</td>
<td>1.958</td>
<td>1.0 ± 1.2 × 10⁻¹⁷</td>
<td>8.7 ± 0.1 × 10⁻¹⁸</td>
<td>0.96 0.93 1.02</td>
<td>0.81 24.0 23.2 0.12</td>
</tr>
<tr>
<td>H₂ Q(1)</td>
<td>2.407</td>
<td>2.0 ± 0.1 × 10⁻¹⁷</td>
<td>2.3 ± 0.8 × 10⁻¹⁷</td>
<td>1.84 2.67 0.7</td>
<td>1.58 0.01 0.01 2.43</td>
</tr>
<tr>
<td>H₂ Q(2)</td>
<td>2.413</td>
<td>5.3 ± 0.3 × 10⁻¹⁸</td>
<td>7.9 ± 0.2 × 10⁻¹⁸</td>
<td>0.49 0.23 0.37</td>
<td>0.005 0.49</td>
</tr>
<tr>
<td>H₂ Q(3)</td>
<td>2.424</td>
<td>3.1 ± 0.1 × 10⁻¹⁷</td>
<td>2.7 ± 0.8 × 10⁻¹⁷</td>
<td>2.79 0.70 0.71</td>
<td>0.006 0.7</td>
</tr>
<tr>
<td>H₂ Q(4)</td>
<td>2.438</td>
<td>1.3 ± 0.1 × 10⁻¹⁷</td>
<td>2.7 ± 0.3 × 10⁻¹⁷</td>
<td>1.2 0.21 0.16</td>
<td>0.002 0.07</td>
</tr>
<tr>
<td>H₂ Q(5)</td>
<td>2.455</td>
<td>1.2 ± 0.1 × 10⁻¹⁷</td>
<td>5.2 ± 0.5 × 10⁻¹⁸</td>
<td>1.11 0.59</td>
<td>0.001 0.05</td>
</tr>
</tbody>
</table>

Columns 5 and 6 list the observed flux for various transitions relative to the S (1) line, for two positions located a) on source and b) ∼ 1″ down the slit. Columns 7 and 8 contain the line ratios expected, in the absence of extinction for thermal excitation at 2000 K (Black & van Dishoek 1987) and for a PDR model taken from Draine & Bertoldi (1996) for n_H = 10⁴ cm⁻³, χ = 10⁵ and T_o = 1000 K.

The line which at best detected at a low significance level corresponds to a blend of the 213/2 - 5/2 transitions at 84.42 and 84.59 μm respectively (see Fig. 3). The integrated flux contained in these two blended lines is 1.51 × 10⁻¹⁰ W cm⁻². Other OH lines within the ISO spectral bands are the 211/2 - 9/2 transitions at 55.9 μm, the 213/2 - 5/2 lines at 119.3 μm and the 213/2 - 1/2 transitions at 163 μm, for which we set 3σ upper limits of 3.1 × 10⁻²⁰ W cm⁻², 2.9 × 10⁻²⁰ W cm⁻² and 5.9 × 10⁻²⁰ W cm⁻² respectively.

### 4.4. H₂O

Water emission has already been observed towards a number of molecular outflows with the ISO spectrometers (Liseau et al. 1996; Ceccarelli et al. 1999; Nisini et al. 1999). It is already well understood that the presence or absence of H₂O in a cloud, may be used to trace the shock activity of the gas (Bergin et al. 1999, 1998). We have been unable to identify any H₂O emission from IRS 5 (see Fig 4), although we set a 2σ upper limit on the 40.69 μm 432–533 line of 3.62 ± 1.55 × 10⁻¹⁵ W m⁻². This is the most intense water line seen in W Hya (Neufeld et al. 1996). No lines were detected towards L1551 IRS 5 at 29.8, 31.8, 174.6, 179.5 and 180.5 μm, which correspond to strong lines that have been detected from other sources, to upper 2σ upper limits of 21, 26, 4, 6 and 6 × 10⁻²⁰ W cm⁻² respectively (uncorrected for extinction – see Fig. 18).

### 4.5. H₂

No emission lines were detected from any of the rotational lines of H₂ in the SWS spectrum towards IRS 5. A single detection of the ν = 1–0 vibrationally excited S (1) line at 2.122 μm has previously been reported by Carr et al. (1987), along with the Q-branch lines at 2.407, 2.414 and 2.424 μm, in a 3″/6 aperture.

![Fig. 22. UKIRT archive spectrum towards L1551 IRS 5. This intensity at the (0,0) position is about three times weaker than the Carr et al. spectrum – but in reasonable agreement since the emission is extended relative to the UKIRT beam.](image)

Their reported S (1) flux of 2.8 ± 0.3 × 10⁻¹⁷ W m⁻² lies below the SWS 2σ sensitivity limit at slightly longer wavelengths (the SWS spectrum starts at 2.4 μm) of 5.6 × 10⁻¹⁶ W m⁻². In view of the low S/N of the Carr et al. (1987) data, we searched the UKIRT archive for a 2 μm spectrum to confirm the Carr et al. result. This spectrum is shown in Fig. 22.

The fluxes in the UKIRT observations are summarised in Table 7. From the upper limits to the ν = 0–0 lines, it is possible to set limits on the beam averaged H₂ column densities using an ‘excitation diagram’. The extinction corrected intensity of an H₂ line I(ν′ J′) is related to the column density of the line, N(ν′ J′) by the relationship:

\[
I(\nu' J') = \frac{hc}{4\pi} \nu A(\nu' J') N(\nu' J')
\]
where $\nu$ are the wave-numbers (Dabrowski 1984) and $A(\nu' J')$ are the transition probabilities for the various transitions, taken from Turner et al. (1977). The column densities are then compared with those predicted for a thermal distribution characterised by a rotational temperature $T_{\text{rot}}$ since the rotational populations of a given vibrational level can be approximated by a thermal distribution:

$$N(\nu' J')/g_j \propto \exp \left( - \frac{E(\nu' J')}{k T_{\text{rot}}} \right)$$

In this relationship the $g$ terms are the degeneracies of the transitions, $g_j = (2J + 1)$ for even $J$ (para-$H_2$) and $g_j = 3(2J + 1)$ for odd $J$ (ortho-$H_2$), and the $E(\nu' J')$ terms represent the energy of the ($\nu' J'$) level. Thus the rotational temperature can be estimated from the inverse of the slope of a plot of $\ln (N(\nu' J')/g_j)$ against $E(\nu' J')/k$, correcting for an ortho/para ratio of 3:1.

It is clear from Fig. 23 that the $\nu = 0$–0 lines are far less affected by extinction than the $\nu = 1$–0 $S(1)$ lines. We can infer from the $2\sigma$ upper limits that the extinction to the $\nu = 1$–0 $S(1)$ emitting gas must be $\lesssim 80^m$ – otherwise we should have detected emission in the $\nu = 0$–0 $S(6)$ or $S(7)$ lines, assuming that the rotational temperature is characteristic of many sources, $T_{\text{rot}} \sim 2000$ K. However, this limit is dominated by the S/N of the UKIRT spectra, and is not usefully stringent alone.

Shock models of the $H_2$ lines have been calculated by a number of workers. For $J$-shocks, 1D models have been presented by Brand et al. (1988), Hollenbach & McKee (1989), Burton et al. (1992), Neufeld & Hollenbach (1994), and for $C$-shocks, by Draine et al. (1983), Smith (1991), Kaufman & Neufeld (1996).

The ISOCAM CVF observations of Cabrit et al. (1999) made towards a number of molecular outflow sources show that the mid-IR $H_2$ lines dominate the cooling, and probe rotational temperatures in the range 300–2000 K. In the cases they studied, the observations were consistent with excitation in low-velocity $C$-shocks ($10$–$30$ km s$^{-1}$), based on the non-detection of $J$-shock tracers predicted to be present in the standard shock models of Hollenbach & McKee (1989), for shock velocities $\gtrsim 50$ km s$^{-1}$. From the L1551 IRS 5 data set, we set $2\sigma$ upper limits of $2.5 \times 10^{-15}$ W m$^{-2}$ and $3.8 \times 10^{-15}$ W m$^{-2}$ on the Ni II $\lambda 6.64$ and Ne II $\lambda 12.81$ lines respectively.

To understand the excitation mechanism of the $H_2$, the $\nu = 2$–1 $S(1) \lambda 2.248$ to $\nu = 1$–0 $S(1) \lambda 2.122$ ratio has often been used as a diagnostic. This ratio has values typically $\sim 0.1$ in shocked regions and molecular outflows, and $\sim 0.6$ for pure fluorescence. However, Draine & Bertoldi (1996) show that in dense PDRs, thermal collisions can transfer the lower-level population ($\nu \leq 2$) towards LTE conditions, so that the line ratios approach those in shocked regions. Observational studies (Usuda et al. 1996; Takami et al. 1999) seem to confirm that this may indeed be the case. In one example, the case of the shocked region close to Orion KL, the ratio varies between $\sim 0.1$ and $0.2$ over a wide range of $\nu = 1$–0 $S(1)$ intensity, whereas in a typical dense PDR such as the Orion Bright Bar, and the reflection nebulae NGC 2023 and NGC 7023, the ratio varies between $\sim 0.2$ and $0.6$, and shows a clear anti-correlation with the $\nu = 1$–0 $S(1)$ intensity. In L1551, the $2\sigma$ upper limit on the $\nu = 2$–1 $S(1) \lambda 2.248$ line is $9.1 \times 10^{-19}$ W m$^{-2}$, and the $\nu = 2$–1 $S(1) \lambda 2.248$ to $\nu = 1$–0 $S(1) \lambda 2.122$ ratio is $\lesssim 0.11$. Thus the ratio appears to be inconsistent with fluorescent excitation, even allowing for thermal collisions. This conclusion is consistent with the lack of an obvious ionising source in the centre of L1551. Although no firm conclusions can be reached on the basis of the line ratio data shown in Table 7, we can rule out (a) low density ($n_{H_2} \sim 10^4$ cm$^{-3}$) shock models and (b) that the $H_2$ emission is seen through more than two or three magnitudes of visual extinction (because the $S(2)$ and $S(3)$ lines, which suffer from the highest extinction, would become too bright relative to any of the models). It thus seems most likely that the $\nu = 1$–0 $H_2$ emission is seen in reflection.

### 4.6 CO vibrational bands

The CO absorption against IRS 5 has been previously studied by Carr et al. (1987) although the new UKIRT data presented here have higher sensitivity and spectral resolution. The data shown in Fig. 24 are fitted by a gas temperature of 2500 K, Doppler width of 5 km s$^{-1}$, and CO gas phase column density of $6 \times 10^{20}$ cm$^{-2}$.
Fig. 24. CO gas modelling and ISO observations of IRS 5. For the upper pane, a model with $T_{\text{ex}} = 2500 \text{ K}$ and Doppler width $\nu_D = 5 \text{ km s}^{-1}$ was run. The best fit column density $= 6 \times 10^{20} \text{ cm}^{-2}$ is shown. Such high temperatures are needed to excite the CO bandhead, and to give the correct relative intensity ratios between the lines. The data and model have a resolution of 500.

4.7. CO rotational lines

It was not possible to detect rotational molecular CO line emission in any of the spectra. For the lowest-$J$ transitions towards L1551 IRS 5, we set 2$\sigma$ upper limits for the $J=14–13$ and $J=15–14$ lines of $4.3 \times 10^{-17} \text{ W m}^{-2}$ and $1.37 \times 10^{-16} \text{ W m}^{-2}$ for L1551 IRS 5 and HH 29 respectively. L1551 appears unlike many other outflow sources, which are rich in shock excited CO line emission.

4.8. O I

Towards L1551 IRS 5 the OI 63um flux is observed to be more intense than along the flow, which indicates that the emission is intrinsic and not due to the diffuse PDR. If the emission was due to the jet, then the mass loss rate is $\sim 7.6 \times 10^{-7} \text{ M}_\odot \text{ yr}^{-1}$ (using the relationship from Hollenbach 1985), which is very similar to the $\sim 10^{-6} \text{ M}_\odot \text{ yr}^{-1}$ derived from this source by other measurements. Along the flow it is likely that the OI emission comes from shocks in the outflows more than from diffuse PDR, since the OI lines are frequently brighter than the C I $\lambda 157 \mu\text{m}$, particularly towards HH29, L1551 NE and on the b1,b2 and b3 positions, which are associated with peaks of C emission in the outflow map by Rainey et al. (1987) and Moriarty-Schieven & Snell (1988).

4.9. HH 29

A deep (2 hour) observation with the LWS was also made centred on the location of HH 29. The baseline subtracted (fitting a low order polynomial to the continuum level) spectrum is shown in Fig. 24, and is devoid of line emission, except for the O I $\lambda 63.2$ and C II $\lambda 157.7$ lines as listed in Table [1].

4.10. L1551 NE

A red 6 $L_\odot$ source has been reported lying close to IRS 5, which has become known as L1551 NE (Emerson et al. 1984). This source has been imaged by Draper et al. (1985) and Campbell et al. (1988). Moriarty-Schieven et al. (1995) first suggested that L1551 NE was the source of a second outflow close to IRS 5, and subsequent observations by Devine et al. (1999) indicated that L1551 NE may be responsible for driving the Herbig-Haro flow HH454. Observations of the continuum towards L1551 NE were also made, and are shown in Fig. 26. This spectrum shows little evidence of any line emission, other than reported in Table [1].

4.11. Other locations along the outflow

Observations at a number of other locations along the molecular outflow, revealed no emission apart from weak C II and O I lines, and will not be further discussed here.
5. Conclusions

Observations have been made towards the two well known infrared sources L1551 IRS 5 and L1551 NE, and at a number of locations in the molecular outflow, using the LWS and SWS spectrometers on the ISO satellite, and several other near-IR telescopes. The present work possibly adds to the complexity by unveiling a plethora of unexpected phenomena, such as e.g. the occurrence of significant extinction at mid and far infrared wavelengths and the existence of dense and hot regions (thermally stable?) over extended scales. The main results of this study are:

1. The ISO LWS spectrum consists of a relatively strong continuum, superposed with a few weak lines of O I, C II and possibly OH. Emission from other species such as CO or H$_2$O was not detected. This might indicate that either the molecules have been destroyed, perhaps in a shock, or that the environment is unable to excite them to emit in the far and near infrared.

2. The ISO SWS spectrum of L1551 IRS 5 contains solid state absorption lines of CO, CO$_2$, H$_2$O, CH$_4$ and CH$_3$OH, which correspond to column densities of $\gtrsim 7.7 \times 10^{18}$ cm$^{-2}$, $5.4 \times 10^{17}$ cm$^{-2}$, $7.6 \times 10^{16}$ cm$^{-2}$ and $2.6 \times 10^{18}$ cm$^{-2}$ respectively.

3. Examination of archival HST NICMOS images reveals a diffuse conical shaped nebulosity, due to scattered light from the central object, with a jet emanating from L1551 IRS 5. It is likely that the emission in this jet-like feature is dominated by [Fe II] lines.

4. The continuum spectral energy distribution has been modelled using a 2D radiative transfer model. The continuum is well fitted for a central source luminosity of 45 $L_{\odot}$, surrounded by a flared disc with an opening angle of 90$^\circ$. The outer parts of the torus extend to a distance of $\sim 3 \times 10^4$ AU, and has a total (gas + dust) mass of $\sim 13 M_{\odot}$. The extinction towards the outflow is estimated to be $A_v \approx 10$ and the mid-plane optical depth to L1551 IRS 5 to be $\approx 120$. This model provides a good fit to the ISO data, as well as the available HST/NICMOS data, mid-IR maps, submm interferometry, and ground-based photometry with a range of different aperture sizes.

5. On the basis of the above model, an extinction curve has been estimated, which shows that the emission at wavelengths shorter than $\sim 2 \mu$m is due to scattered light from close to L1551 IRS 5, while at wavelengths $\gtrsim 4 \mu$m, is seen through the full extending column towards the central source. This need to be taken careful account of when comparing line intensities at different wavelengths.

6. Three [Fe II] lines were detected in the SWS spectrum towards L1551 IRS 5, with a fourth line at $\sim 1.5 \sigma$, and upper limits on several others. Although it would seem at first sight that shocks would be the most likely source of excitation in a known shocked region such as this, the line intensities do not fit the predictions of simple shock models. The problem with such shock interpretation of the [Fe II] lines is that the line ratios and strengths imply densities and temperatures which are outside the range considered by Hollenbach & McKee. However, as explained in 4.1.1 4.1.3, the overall energy budget and observed densities place stringent limits on possible shock models, leading us to explore other models. An alternative explanation has been examined where the [Fe II] gas is hot ($\sim 4000$ K) and dense ($n_{H_2} \gtrsim 10^9$ cm$^{-3}$). Although this provides an acceptable fit to the relative line intensities, it provides no constraints as to the precise heating mechanism of the gas – although it seems likely that it would have to occur very close to the root of the outflow. The lack of detection on Br $\alpha$ $\lambda 4.052$ or Br $\gamma$ $\lambda 2.626$, and the known low surface temperature ($T_{surf} \sim 5500$ K) of the central protostellar object argue against efficient excitation in a high UV field environment, making shocks the most likely way to explain the [Fe II] and [Si II] emission intensities.

7. The SWS observations did not detect any emission from rotationally excited H$_2$. Observations with UKIRT of the vibrationally excited S and Q-branch lines were consistent with the gas having an excitation temperature of $\sim 2500$ K. Given the likely opacity to the central source which was predicted in our modelling, it is unlikely that we would have detected emission due to this hot gas component. Similarly, there was no evidence of lower temperature ($\sim 500$ K) gas, as has been inferred towards many other sources.

8. Observations with UKIRT of the CO absorption bands close to 2.4 $\mu$m are best fit with gas temperatures $\sim 2500$ K, and a column density $\sim 6 \times 10^{20}$ cm$^{-2}$.

9. Evidence for dense (coronal and higher densities) and hot (at least 2500 K up to perhaps 5000 K) gas is manifest in a multitude of observables; a) the overall SED, b) the CO bands, c) for a ‘normal’ [CO]/[H$_2$] = 8 $10^{-5}$, the implied H$_2$ column density N(H$_2$) = 6 $10^{20}$/8 $10^{-5}$ = 8 $10^{24}$ $\sim 10^{25}$ cm$^{-2}$ for gas at T $\gtrsim 2500$ K, d) the H$_2$ emission (presence of rovibrational emission and absence of pure rotational lines) e) the [Fe II] spectrum, and f) that the SiII 35 micron is also consistent with this gas.

Although each of these individual pieces of evidence may have several interpretations, we argue that the combination of them all makes it highly likely that such a hot and dense gas phase is present. The problem as how to produce and to maintain this region (which has not previously been presented in the literature elsewhere) is beyond the scope of the present paper. Our simple modelling is based on steady state assumptions about processes which are most probably highly dynamic in nature. Any attempt to provide an explanation(s) would at the moment be based on pure speculation, and so we instead indicate the need for extensive theoretical work, which is beyond the scope of the present observational paper.

10. Observations at a number of other locations along the molecular outflow, and towards the Herbig-Haro object HH29, revealed no emission apart from weak C II and O I lines.

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